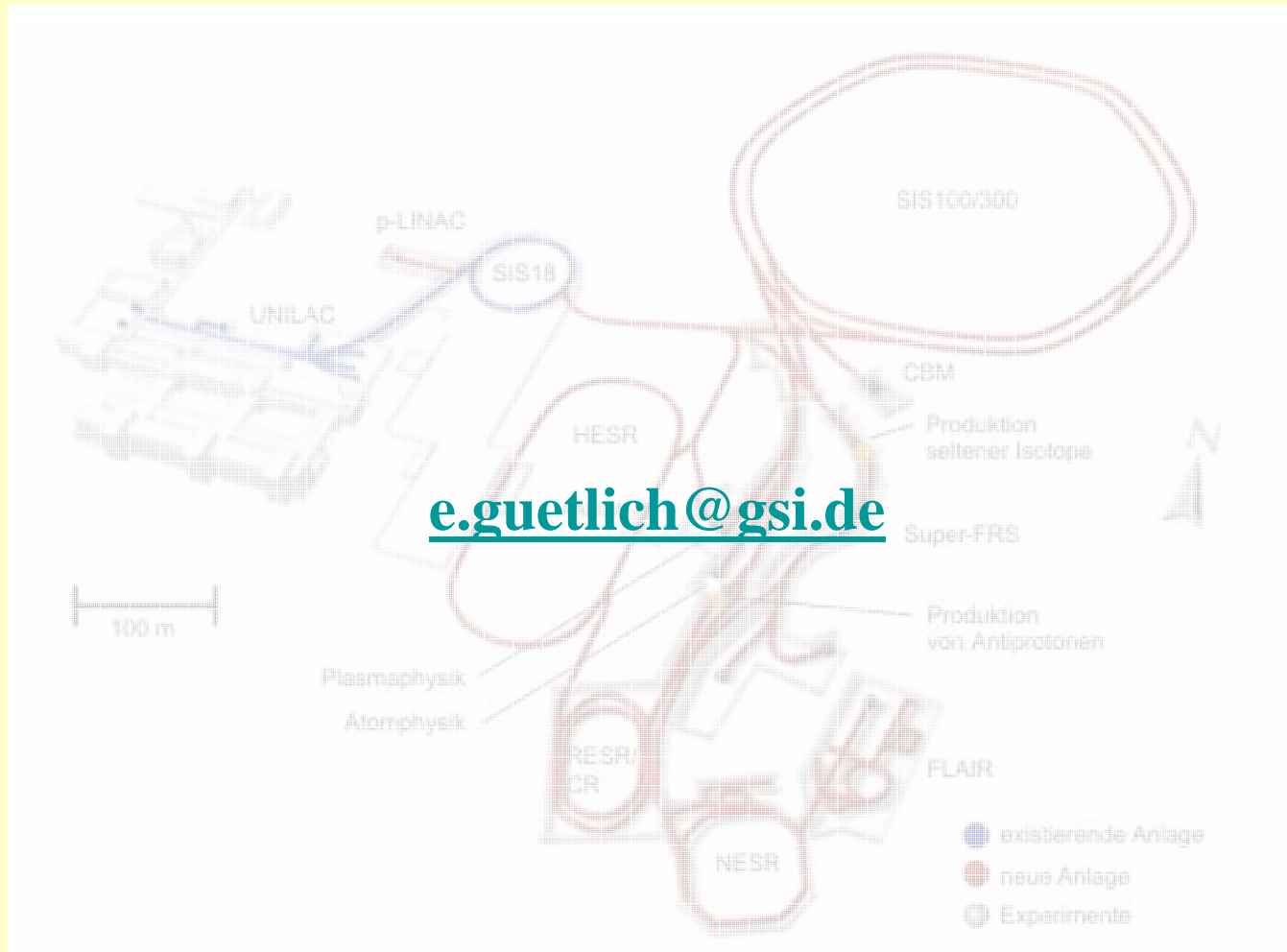


2010, a scintillating year

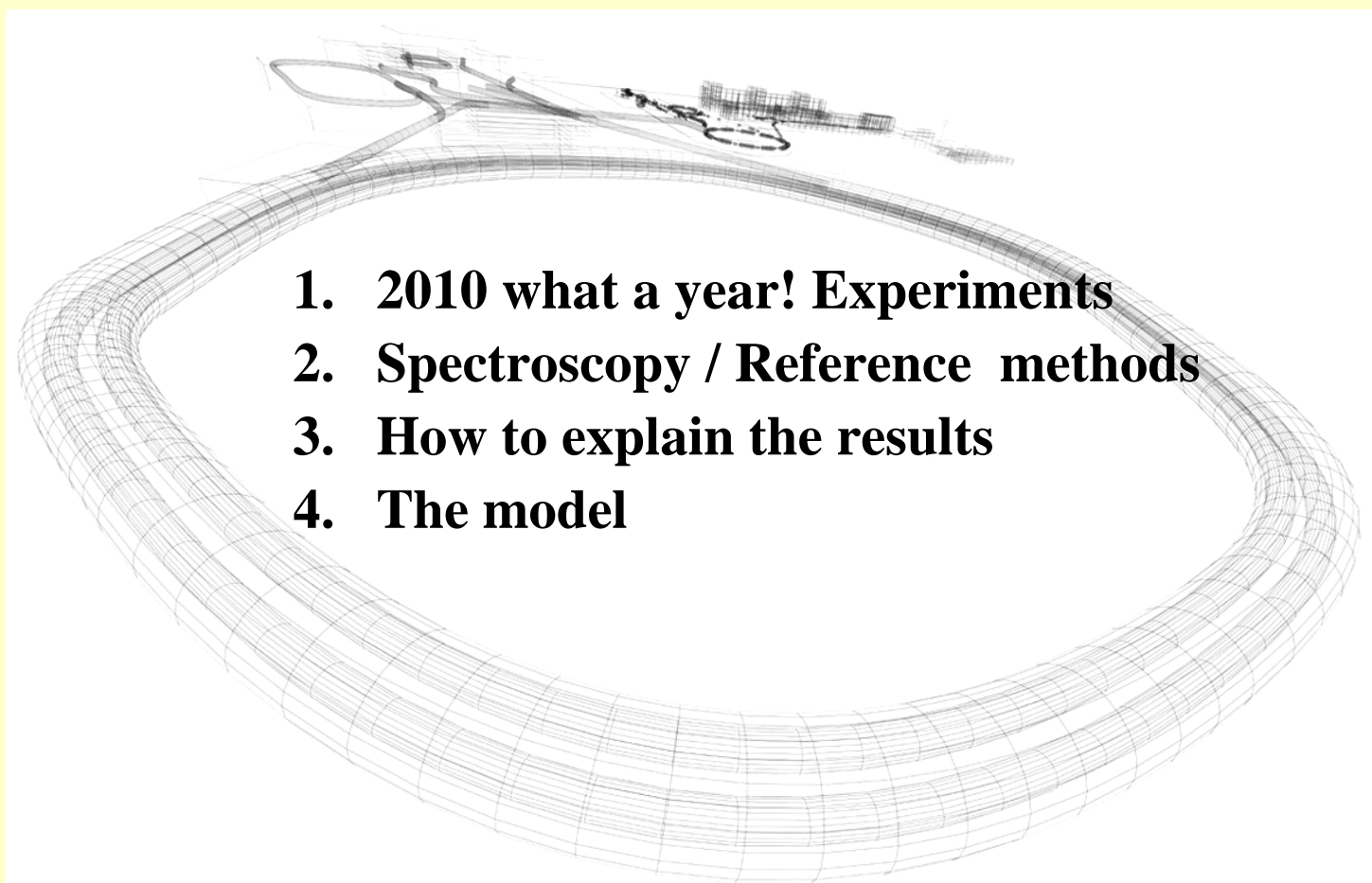
Eiko Gütlich^{1,2}, Peter Forck¹, Wolfgang Ensinger², Beata Walasek-Höhne¹

¹GSI, Darmstadt, Germany

²Technical University Darmstadt, Germany



2010, a scintillating year

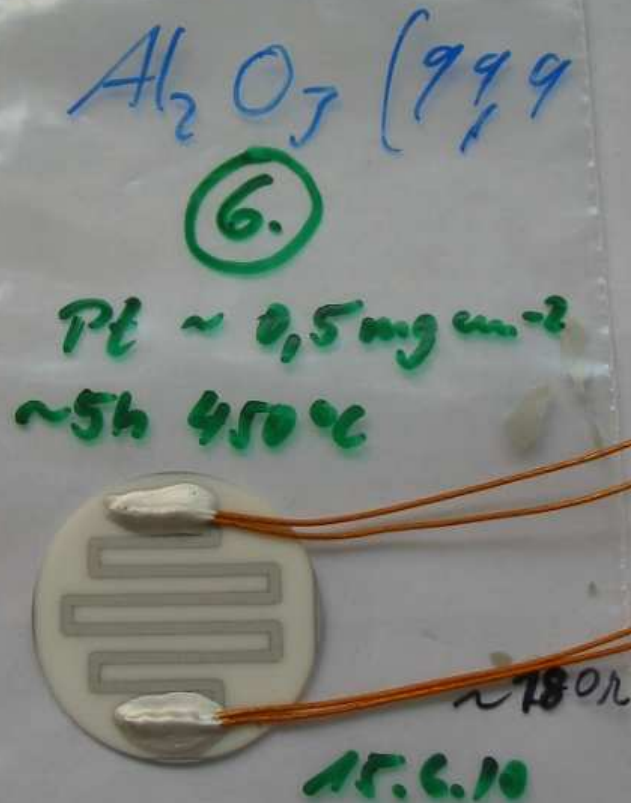


- 1. 2010 what a year! Experiments**
- 2. Spectroscopy / Reference methods**
- 3. How to explain the results**
- 4. The model**

The year 2010

- New experimental design for Spectroscopy and the camera setup
- New heating method
- Beam time marathon: March, May, July → each 8 weeks → lots of data
- Due to the huge amount of new data → a lot of programming in SciLab
- Results up to April have been presented on BIW'10 and DPG-Tagung'10 (Regensburg)
- A lot of literature inquest
- Development of a model
-
-

New heating method



- A Pt layer of 250 nm is sputtered on the backside of the sample.
- The layer is connected to the Capton-wire via high temperature conductive glue Elecolit 327.
- The layer is annealed before characterisation to ensure stable conditions.

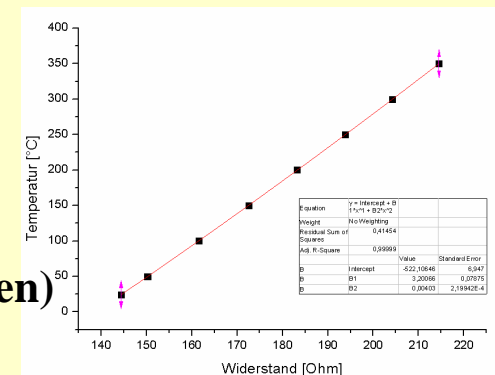
Advantage:

The temperature behavior can be investigated by simultaneous heating and direct 4-point temperature measurement up to **400°C**

350°C

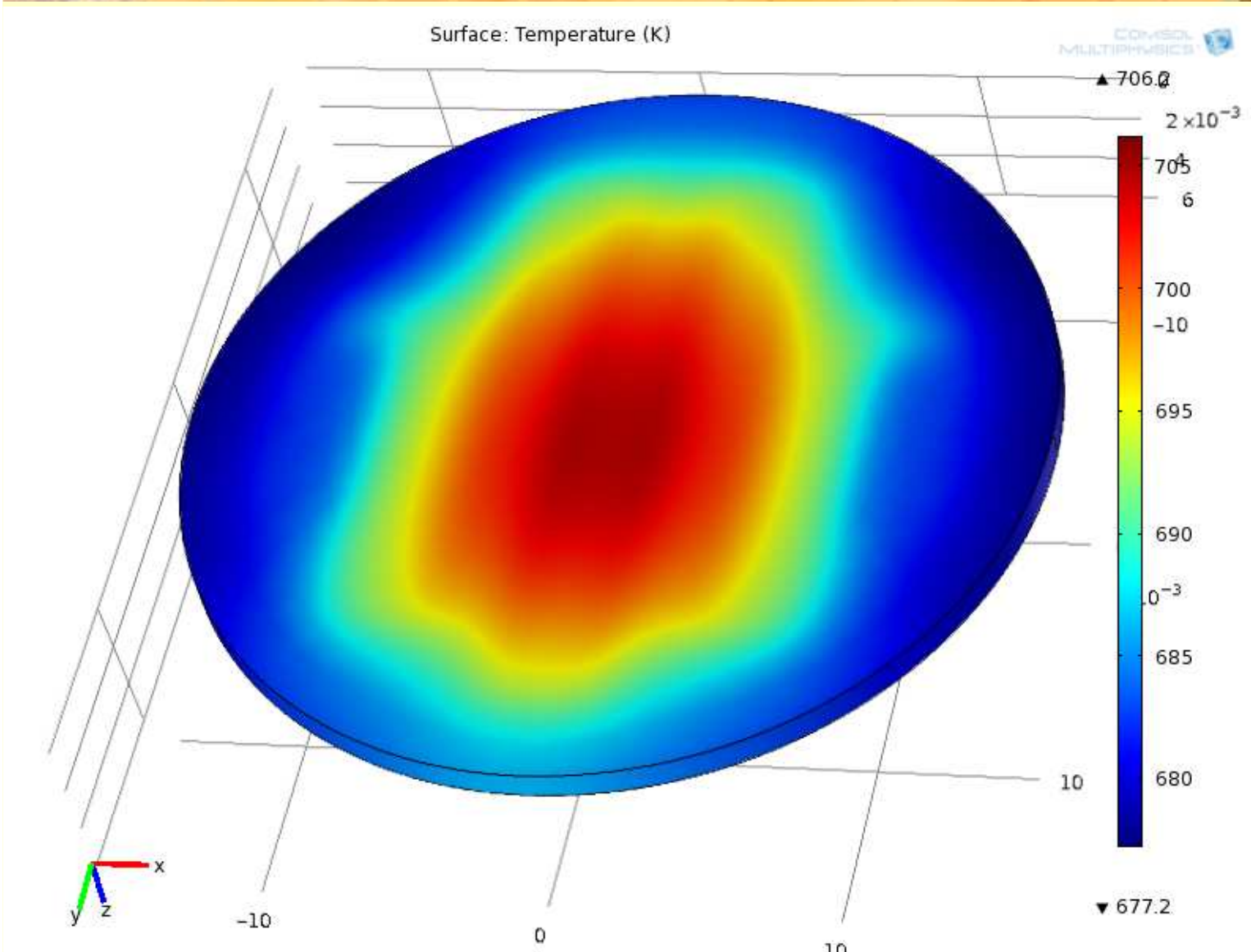
→ 700 A/mm² (Pt)

→ 1,1 W/cm² (Screen)



GS

Simulation of the heating



Results:

- Simulation fit the experimental data
- Temperature difference on the area of the typical beam is always smaller than 10°C (typically 5°C)

Why is the spectrum important ?

The spectral response can give information about the scintillation mechanism. → once you know it, one can try to make predictions

When the spectrum is different between the outer-part and the centre, it could lead to a wrong representation of the ion beam due to:

- **different response of the states to the deposit energy.**
- **the wavelength dependent sensitivity of the CCD-Chip.**
- **the different chromatic aberration (Farblängsfehler) of the lens-system used.**

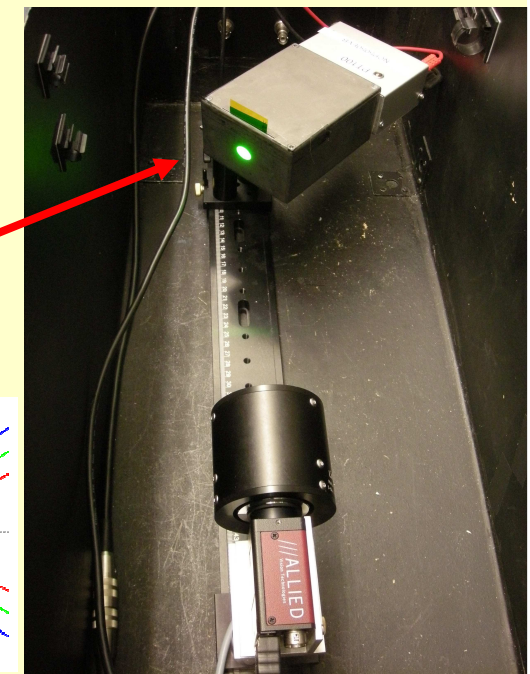
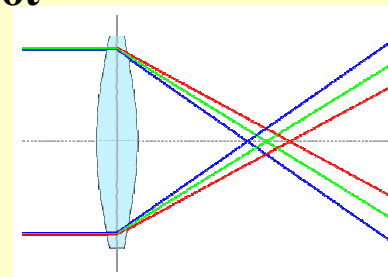
Up to now, the chromatic aberrations (Farblängsfehler) of the used lens systems have been investigated in the 400-800 nm region, with a purpose-build light source "beam-spot simulator". → Error in Sigma within 1%

Linearity of the chip (double integration time → double Pixel value)





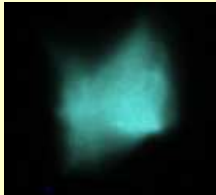
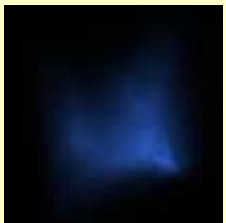
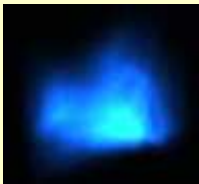

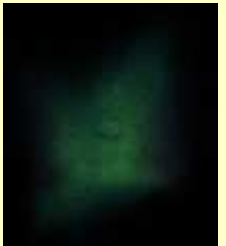
Iris values (Blendenzahlen)



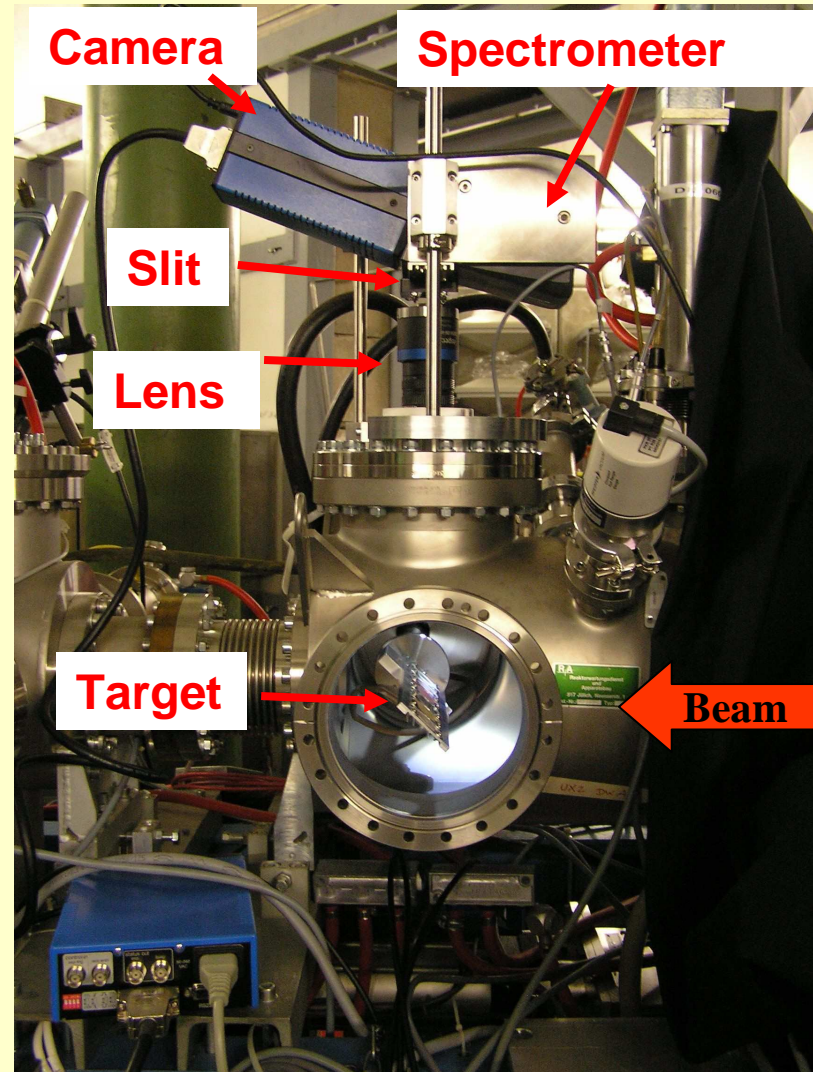
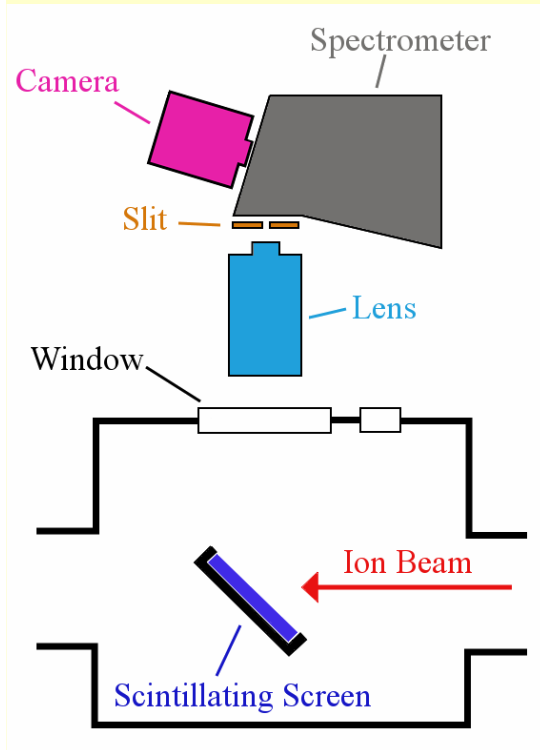
Courtesy of Jan Mäder (GSI)



Measurement with colour CCD

| | $\text{Al}_2\text{O}_3 : \text{Cr}$ | Al_2O_3 | Herasil | ZrO_2 | $\text{ZrO}_2 : \text{Mg}$ | BN |
|-------------------------|---|---|---|---|---|---|
| Within the Pulse |  |  |  |  |  |  |
| Afterglow | the same colour $t_2: \sim 2\text{ms}$ | the same colour $t_2 < 30\mu\text{s}$ |  $t_2 < 30\mu\text{s}$ |  $t_2 < 30\mu\text{s}$ | the same colour $t_2 < 30\mu\text{s}$ |  $t_2 < 30\mu\text{s}$ |

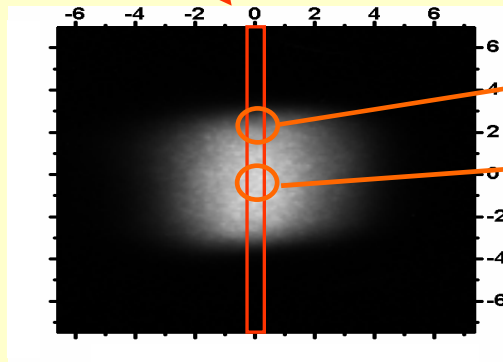
Advanced experimental setup (spectroscopy)



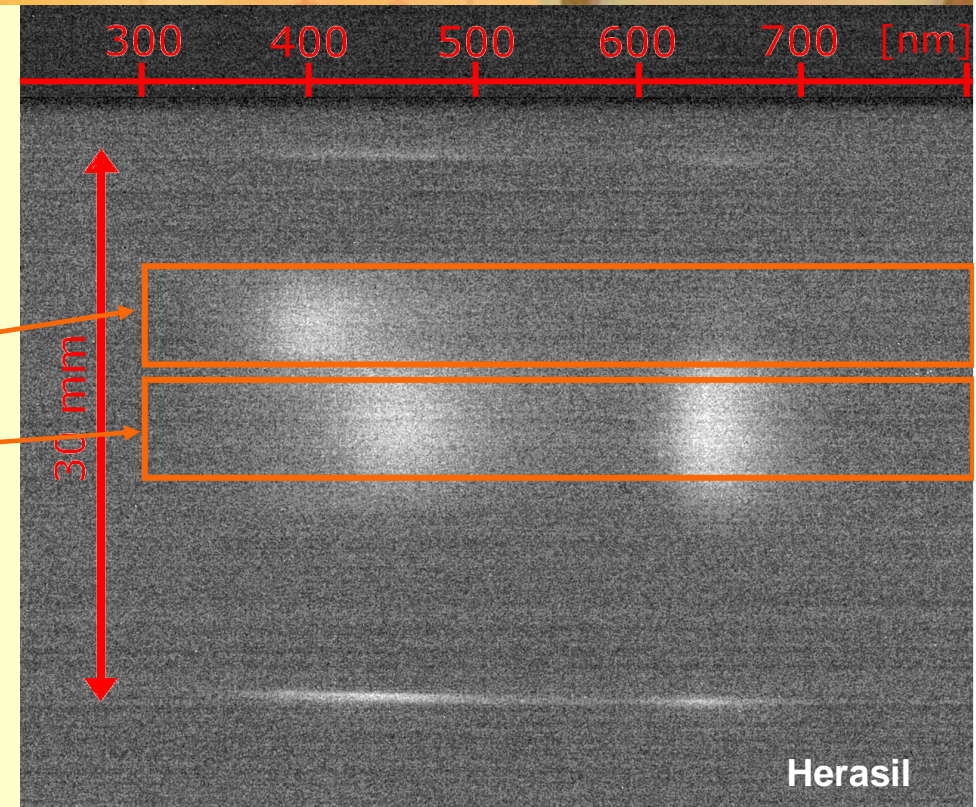
Camera: PCO1600
Spectrometer: Horiba
Jobin Yvon CP140-202
Slit: Newport M-SV-0.5
Lens: Linos inspec.x UV-
Vis-Lens, 50mm focal
length

Spectroscopic investigations

Investigated area of the beam spot for spectroscopic studies



typical width of the stripe is about 1mm and depend on the slit setting

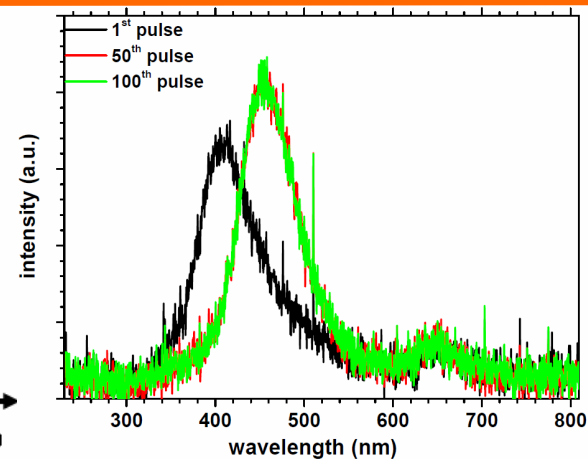
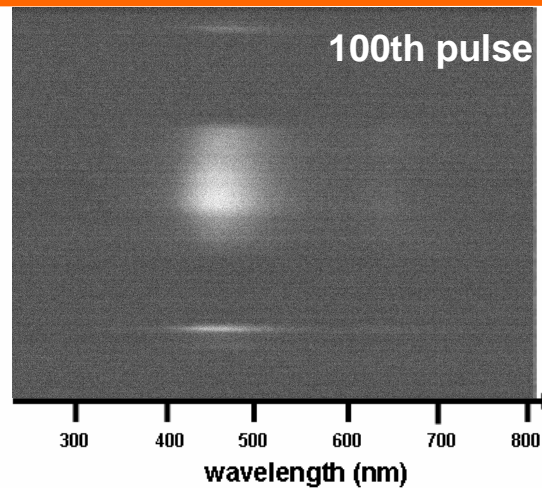
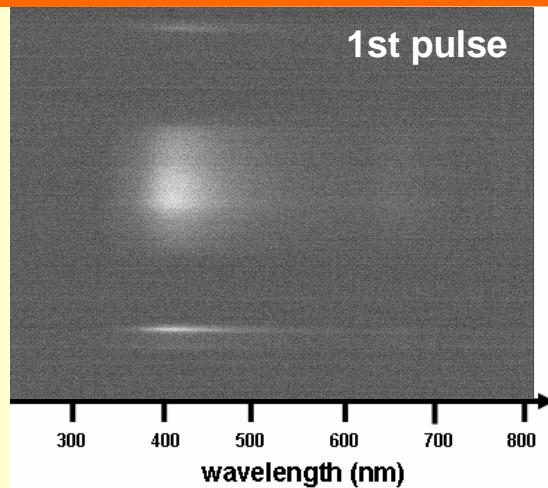
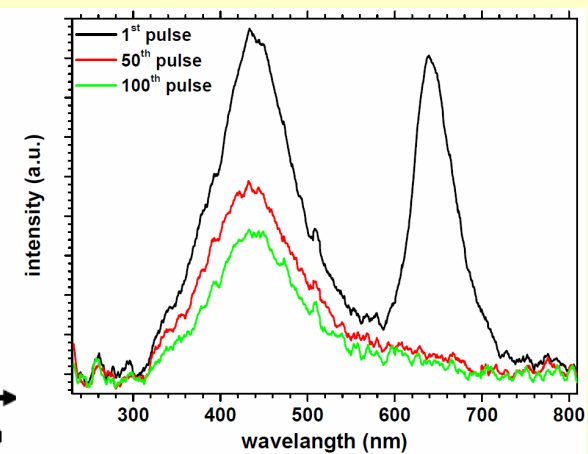
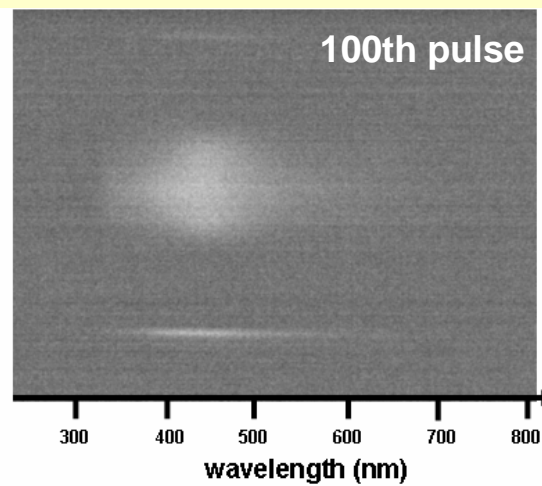
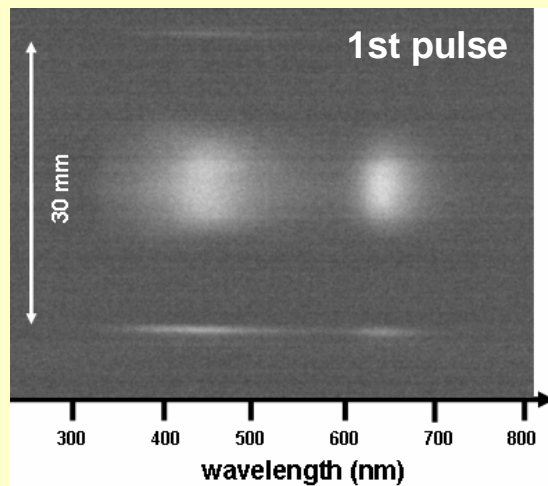


Advantage:

- influence of the ion flux on the spectra can be analysed over the flounce, for each macro pulse
- the whole screen is observed

Spectroscopic investigations on Herasil

Beam parameters: U^{28+} , 4.8MeV/u, $5.2 \cdot 10^{10}$ Ions/Pulse in 0.8ms

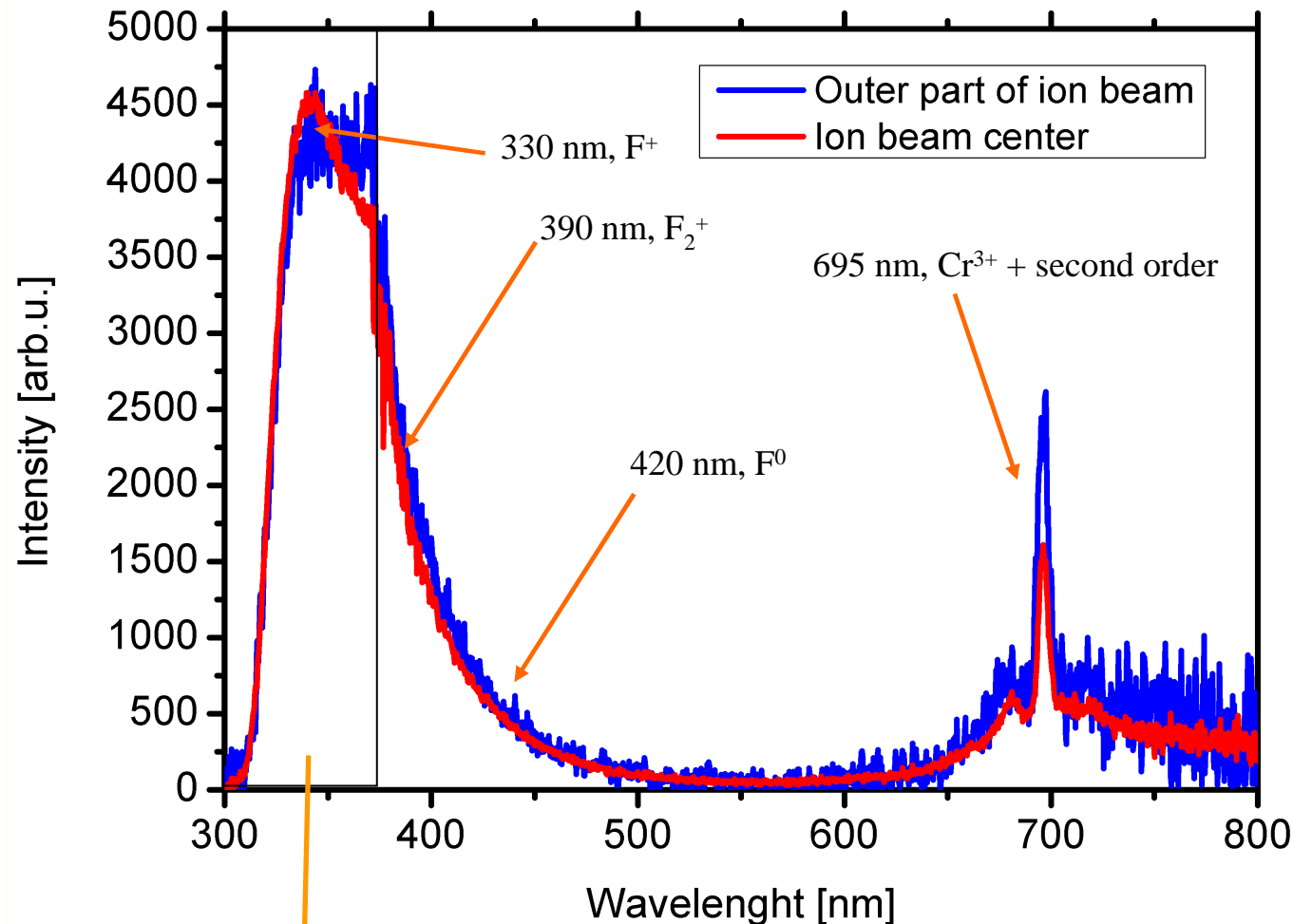


Beam parameters: Ca^{10+} , 4.8MeV/u, $9 \cdot 10^{10}$ Ions/Pulse in 5.3ms

Result: The Spectrum can depend on the ion and on the dose!

Spectrum of Al_2O_3 for $^{48}\text{Ca}^{28+}$ @ 4.8 MeV/u

Beam parameters: $^{48}\text{Ca}^{10+}$, 4.8 MeV/u, $5.2 \cdot 10^{10}$ Ions/Pulse in 3ms, 1Hz



Results:

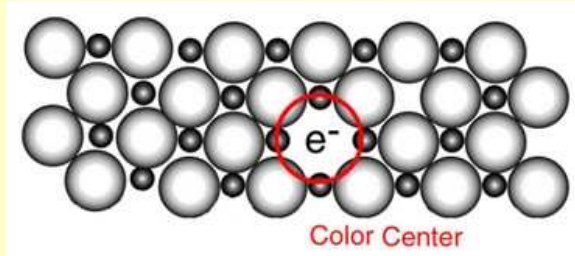
- Spectra of outer and centre are similar to each other
- No $\text{F}^0 \leftrightarrow \text{F}^+$ conversion @ 4.8 and 11.4 MeV/u
- High current results correspond to low current results in the literature

Effective response of the optical system for standard measurements @ 370nm is < 10 % of 500 nm

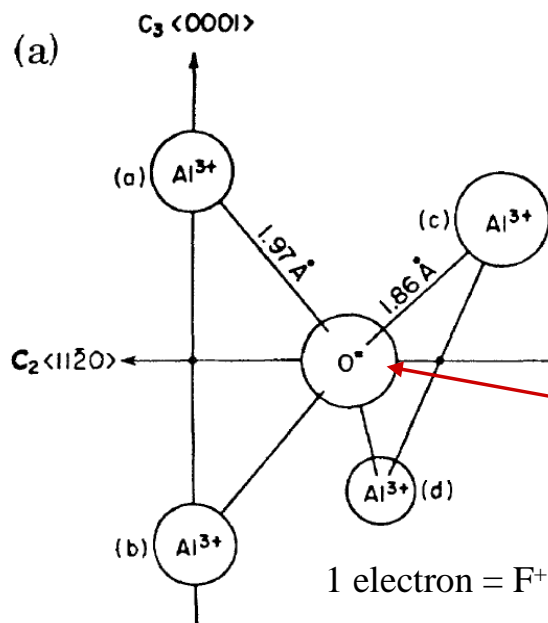
Region that can not be observed with standard, not UV enhanced optics

Colour Centres in Al_2O_3

Example for a color center →



The Al_2O_3 lattice



1 or 2 Electrons



1 electron = F^+ centre, 3.8 eV (326 nm), $\tau=1.7$ ns

2 electrons = F^0 centre, 3.0 eV (413 nm), $\tau=25$ ms

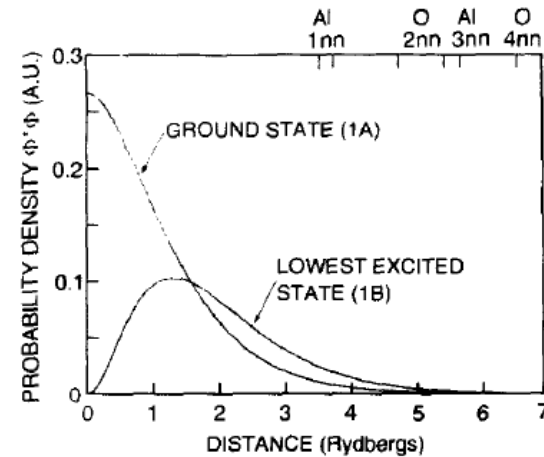
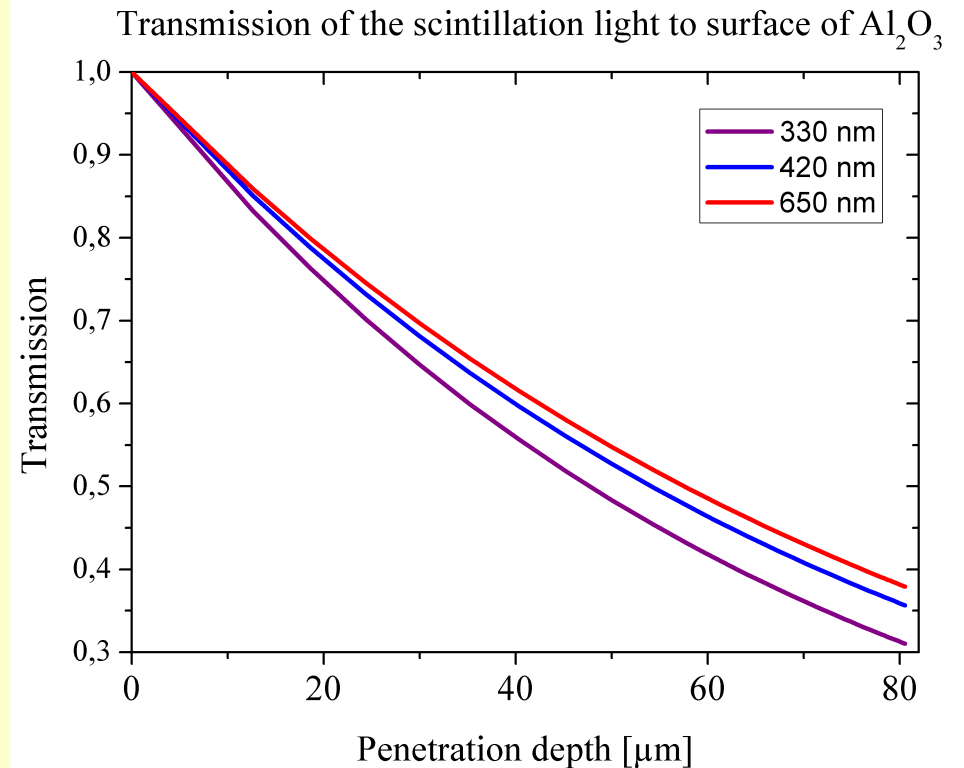
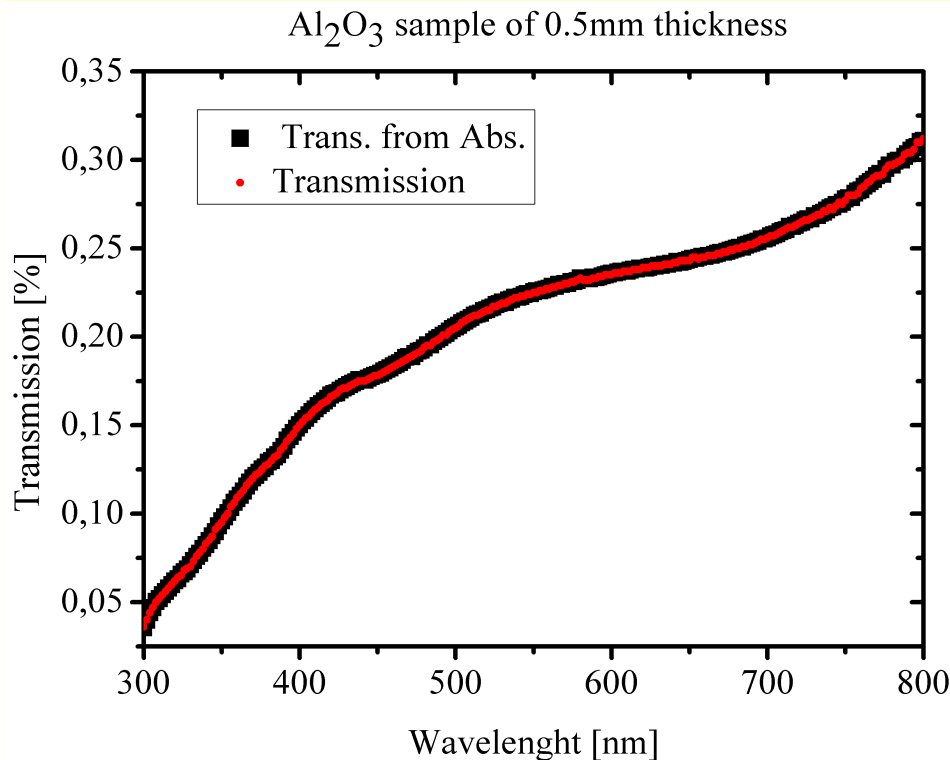


Fig. 2. F^+ center relative probability density $\phi^* \phi$ in $\alpha\text{-Al}_2\text{O}_3$, calculated from wave functions derived in Ref. [38]. For a typical ground state (1A) distribution, that along the y-axis in Fig. 1 is depicted; the 1B-level distribution shown is along the x-axis, out of the plane of Fig. 1.

B.D. Evans / Journal of Nuclear Materials 219 (1995) 202–223

Result: The F^+ emission is might be more resisted against quenching because of the less extended wave function and the shorter live time.

Transmission of Al_2O_3



Result: There is no important contribution from reflections in the material to the detected scintillation signal. → backside to surface reflection @ 420nm would be reduced to $1\text{E}-5$, and for 0.5mm sample to $2\text{E}-11$ for 1mm.

Interesting previous low current investigation

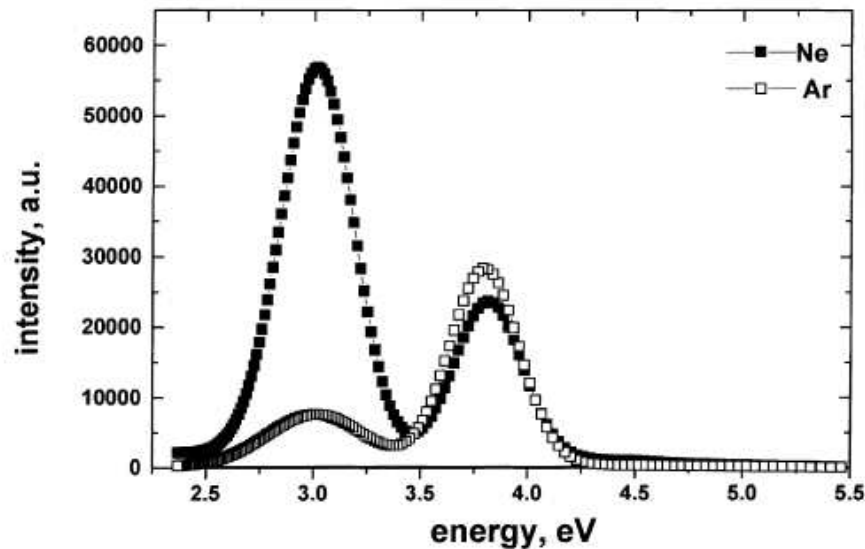


Fig. 5. Luminescence spectra from $\alpha\text{-Al}_2\text{O}_3$ measured under Ne and Ar ion irradiation at $T=80$ K and damage dose $> 10^{-4}$ dpa (intensity saturation stage).

V.A. Skuratov / Nucl. Instr. and Meth. in Phys. Res. B 146 (1998) 385–392

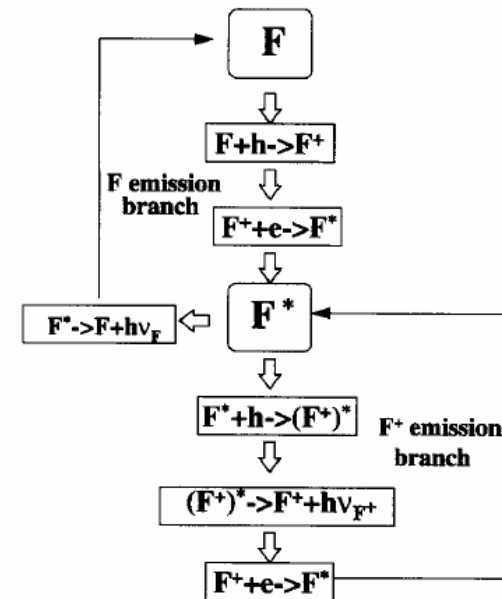


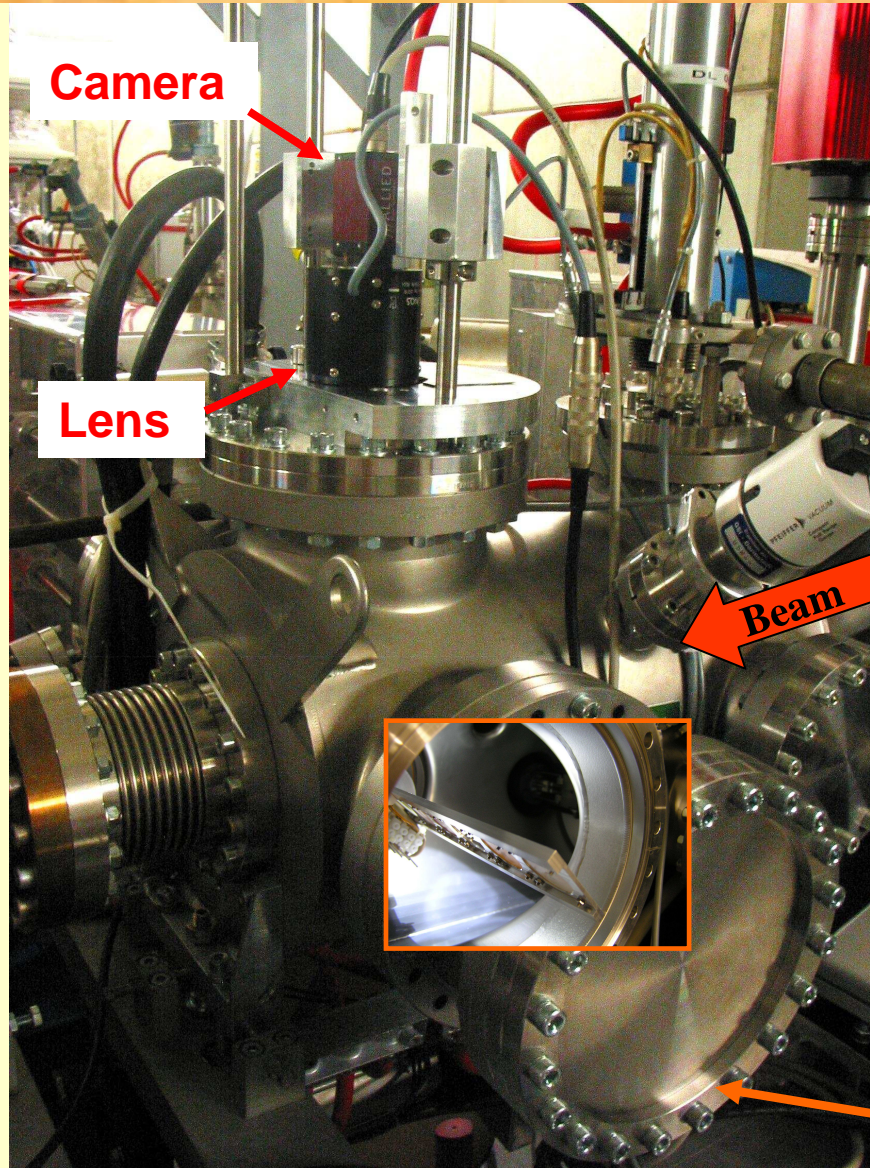
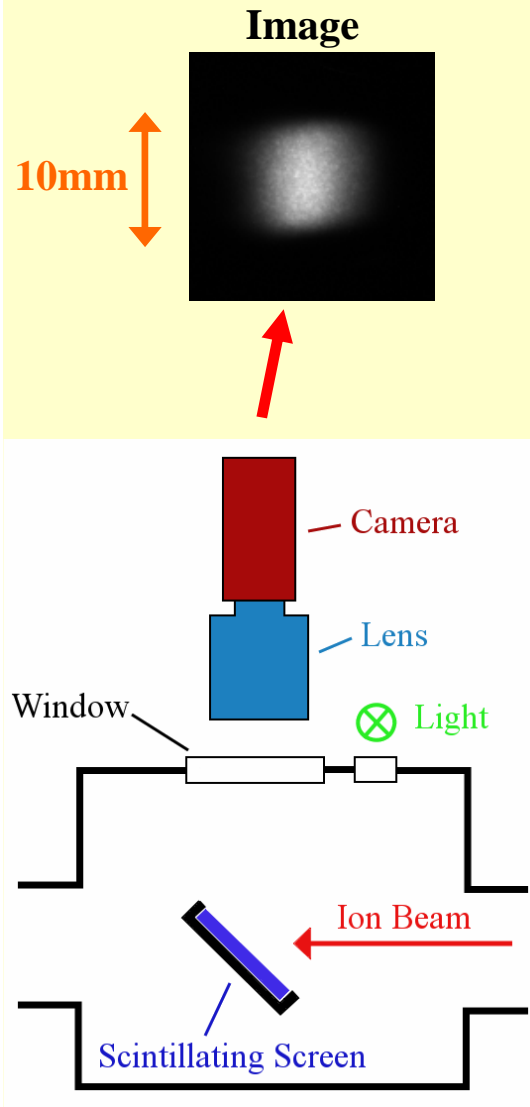
Fig. 6. Flow diagram for the F and F^+ luminescence processes.

A. Morono, E.R. Hodgson / Journal of Nuclear Materials 249 (1997) 128–132

Results: The spectrum of Al_2O_3 depend on the ion species.

It is considered possible that F^+ and F^0 can convert into one another, depending on the ionization desity, i.e. the e-h population .

Advanced experimental setup (imaging prop.)



Camera: AVT Stingray F033B
(VGA monochromatic), FireWire interface
Lens: Linos ROD Mevis, 2516, stepping motor driven
Resolution: 10 pixel/mm
DAQ: Industrial PC with FPGA

Advantage:

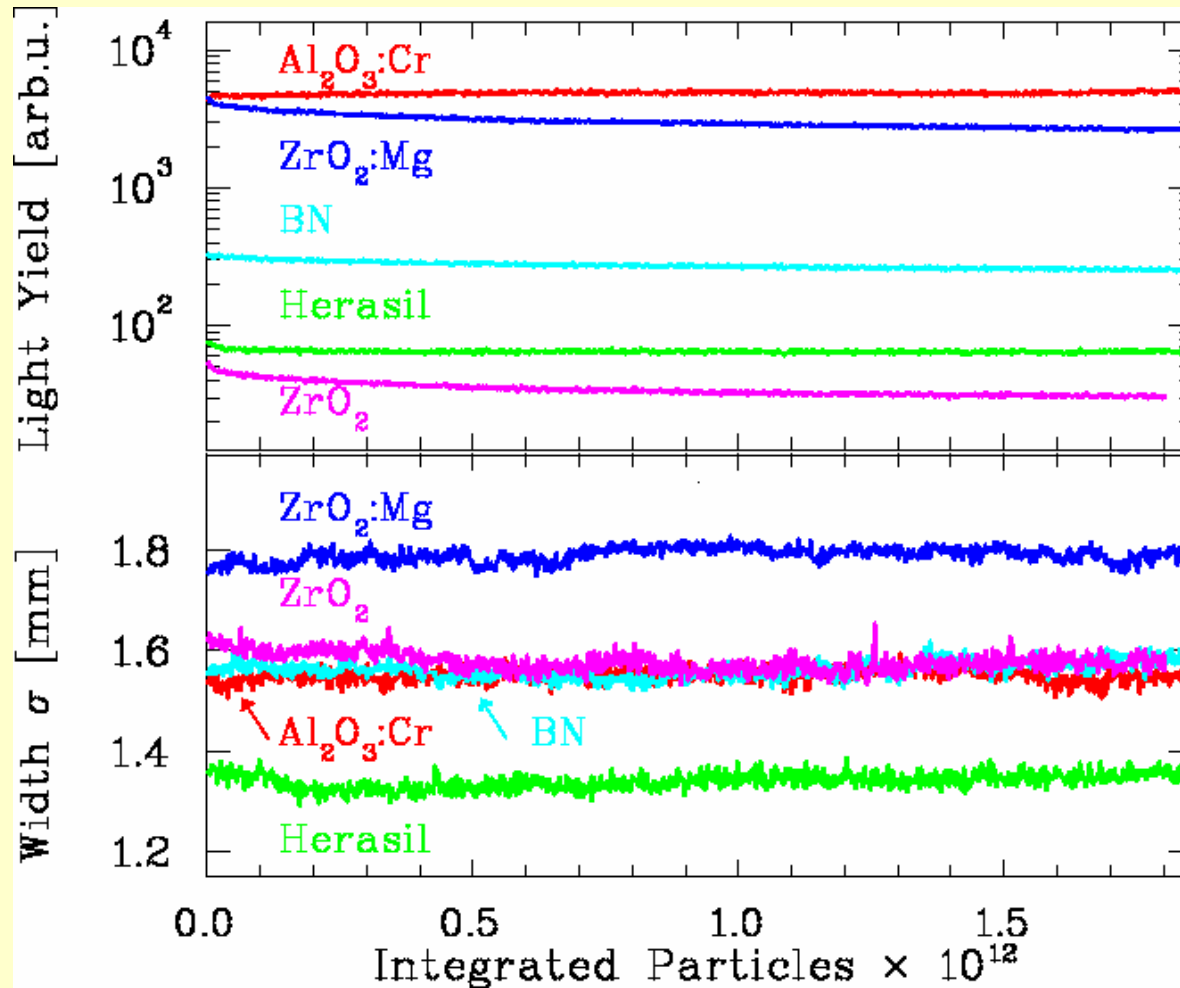
- back-fitting time from spectrometer to normal camera is about ~25 min.
- new DAQ stores the number of particles synchr. for each image → **new@GSI**

Flange diameter 200 mm



Light yield and profile width @ low intensity

Beam parameters: $^{40}\text{Ar}^{10+}$, 11.4 MeV/u, $2 \cdot 10^9$ Ions/Pulse in 100 μs , $\sim 30 \mu\text{A}$, 2.4 Hz, 1000 beam pulses



Results:

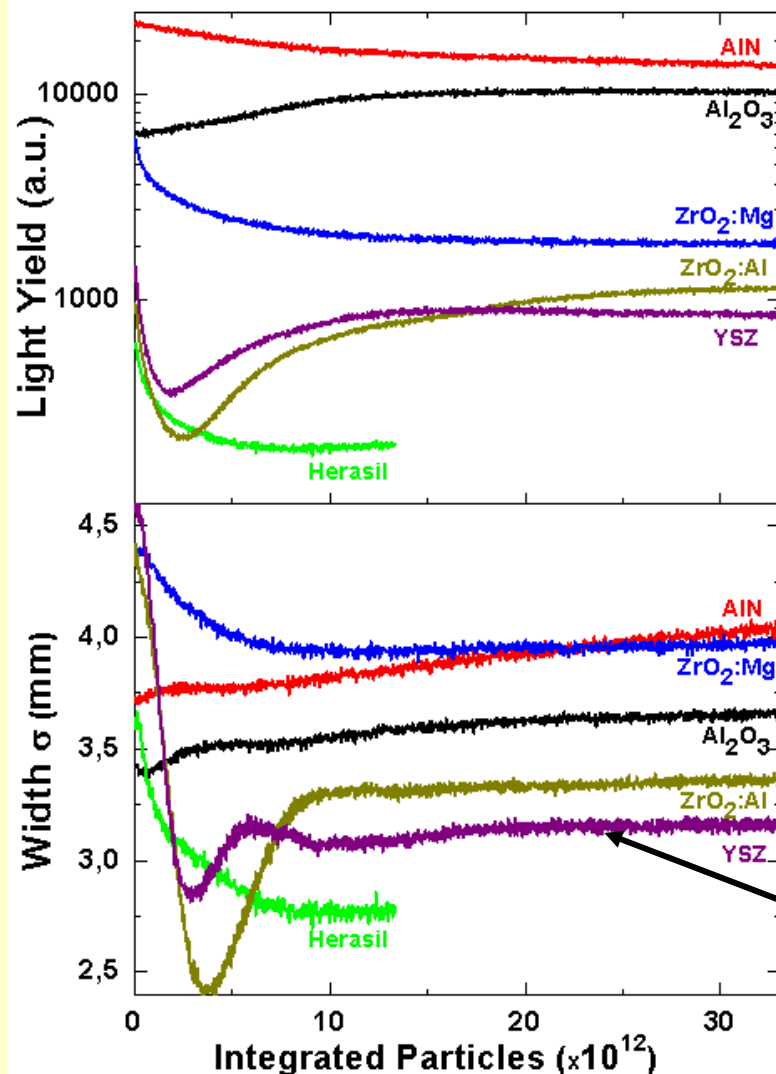
- reproducible behavior
- different light yield and width reading
- light yield does not correlate with beam width
- different beam shape (from higher stat. moments)

Difference of 14% in profile width is not negligible for quantitative evaluation

Average temperature: $\sim 47^\circ\text{C}$
(backside of $\text{ZrO}_2:\text{Mg}$)

Light yield and profile width @ higher intensity

Beam parameters: Ar^{10+} , 11.4 MeV/u, $3.3 \cdot 10^{10}$ Ions/Pulse in 0.2ms, 260 μA , 1.7Hz, 1000 Pulse



~10 times higher beam current

Results:

- light yield and profile width depend on material
- different dynamical behavior
- possible reasons: material modification and temperature dependency
- for the zircon oxides the behavior is clearly temperature dominated

Difference of 30%
in profile width
is not negligible for
quantitative evaluation

Average temperature: $\sim 200^\circ\text{C}$
(backside of Al_2O_3)

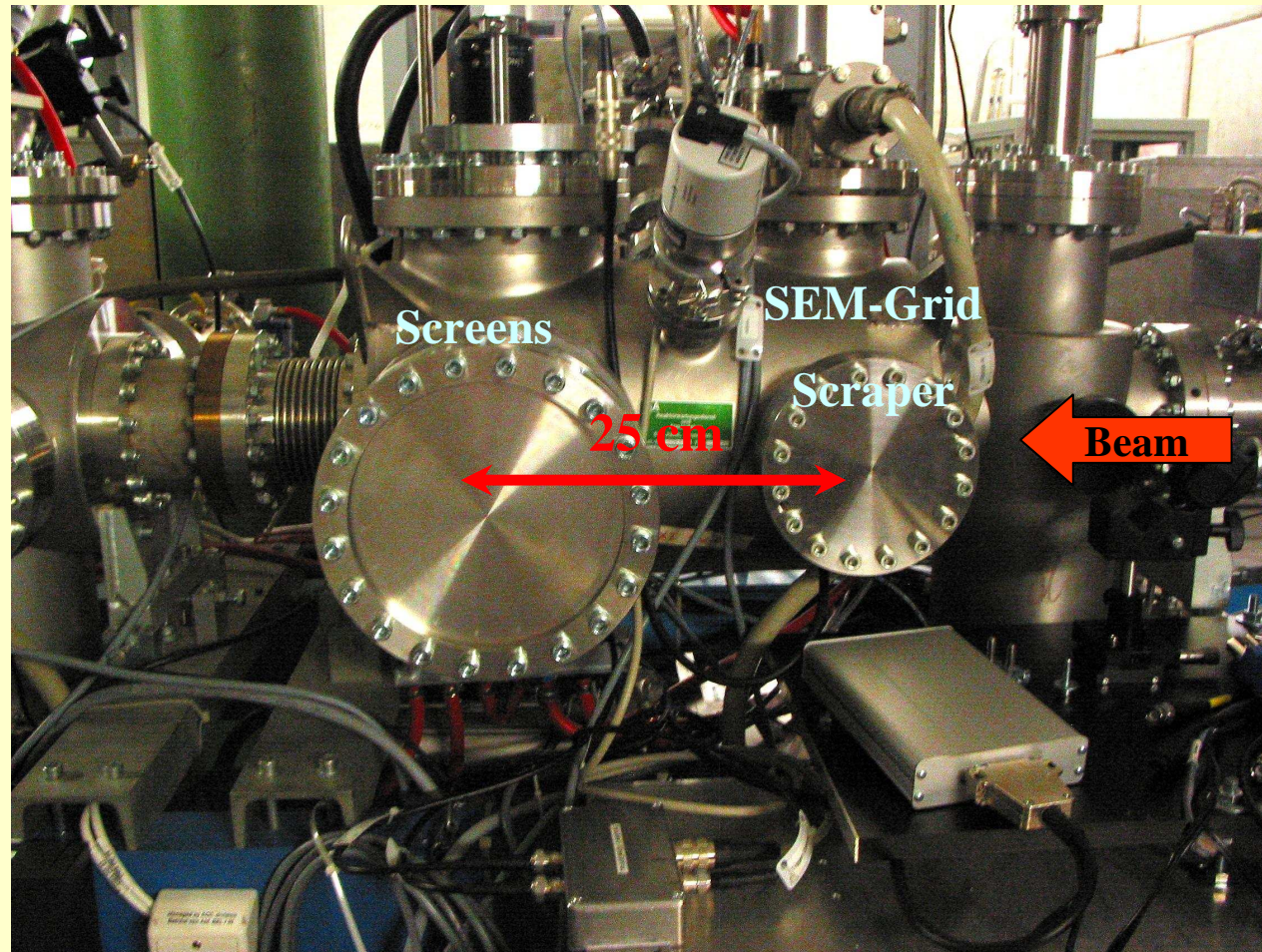
What we have seen up to now

The different materials measure different values for the transversal beam parameters for the same ion beam!

Which one is right, or are they all wrong?

What are the parameters of the ,real‘ ion beam ?

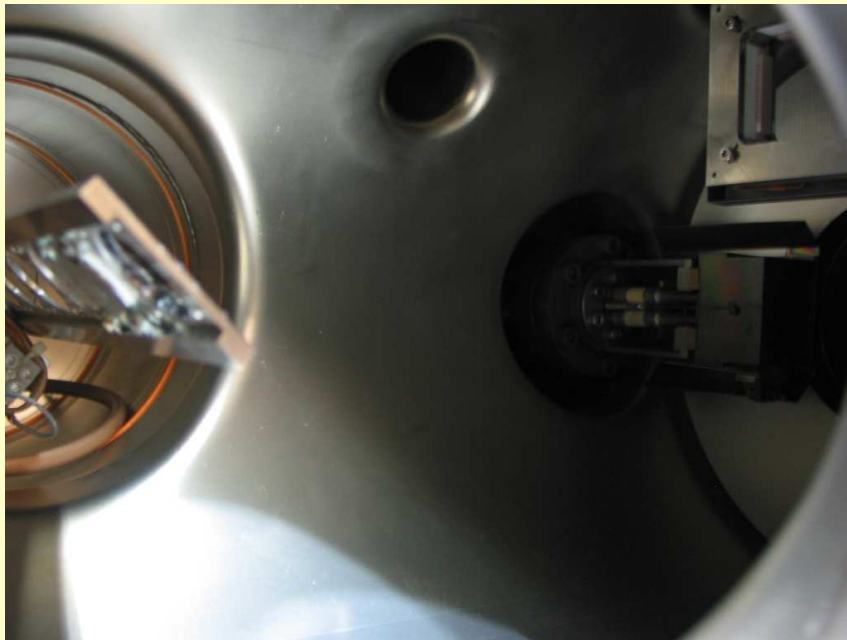
Comparison with reference methods



- One can measure a reference profile 25cm in front of the screens.

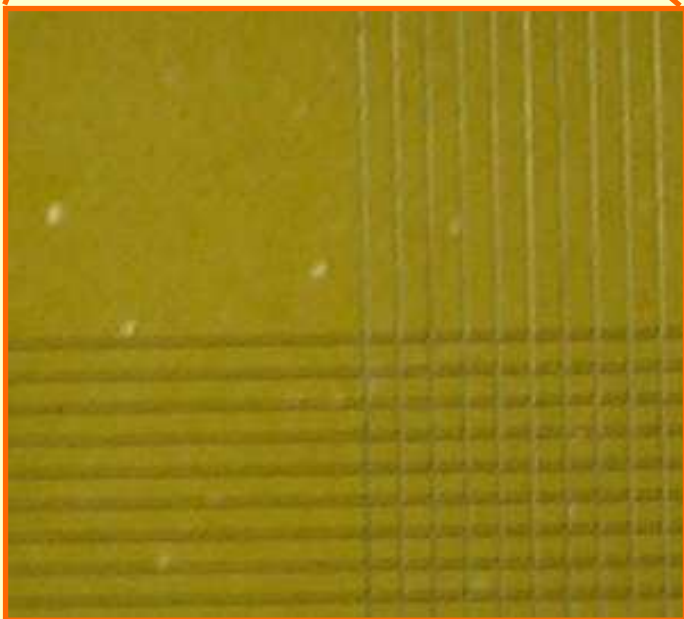
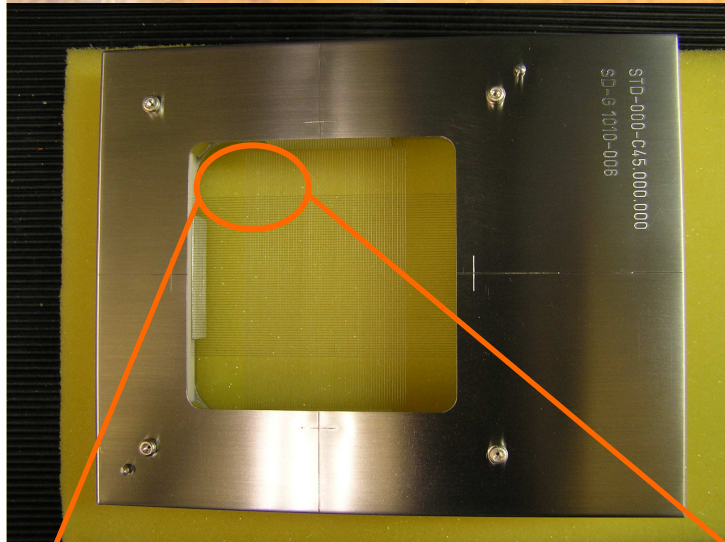
Limitation:

- Due to the lack of space it is not possible to take reference profiles at the same optical position as the screens.
- The profile grid is unable to measure the profile of the entire macro pulse



Secondary electron emission (SEM) grid

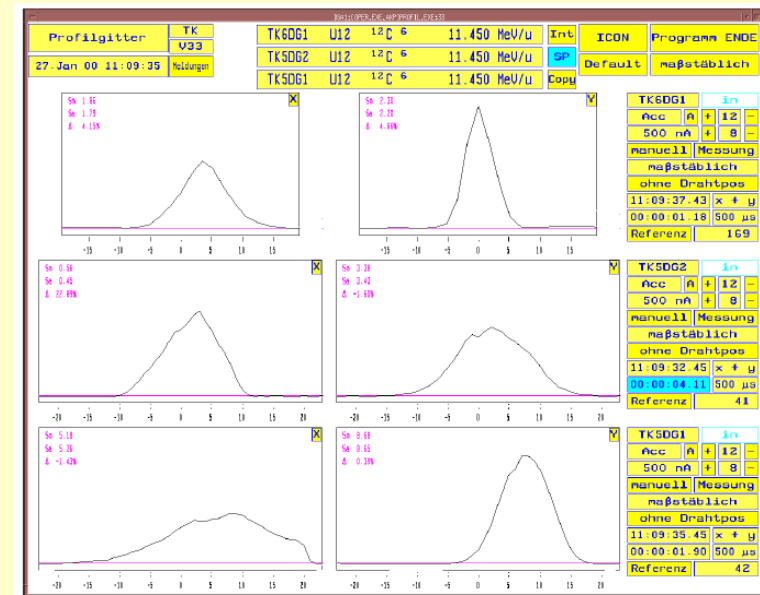
(Standard method)



When particles hit a surface, secondary electrons are liberated and escaping from the surface. For the profile determination, individual wires or ribbons interact with the beam; this is called a Secondary Electron **E**Mission grid.

Limitation:

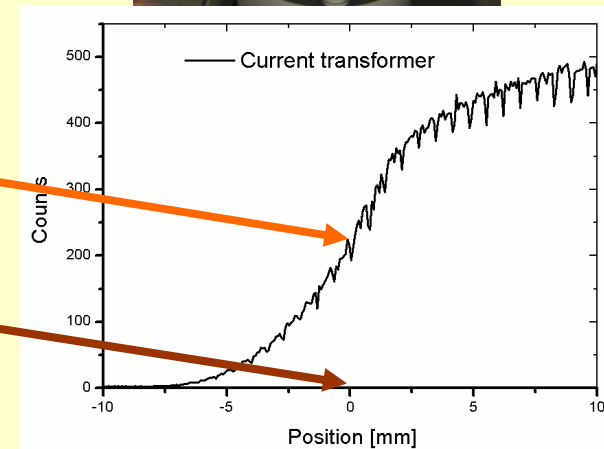
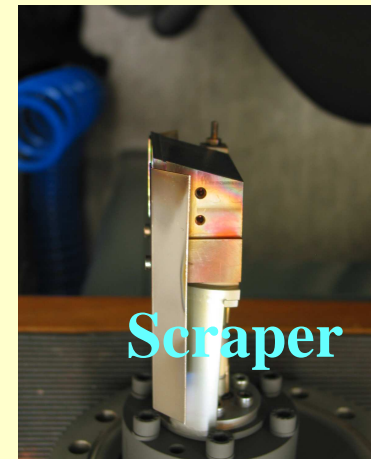
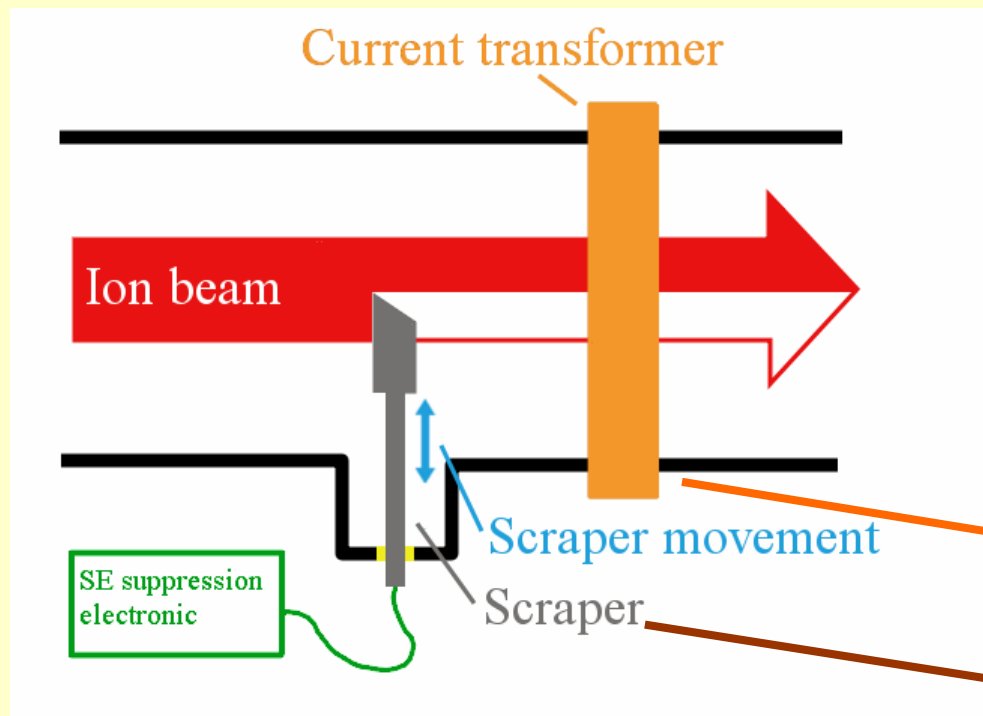
distance between wires:
~ 0.8 - 1.5mm



The 2nd reference method?

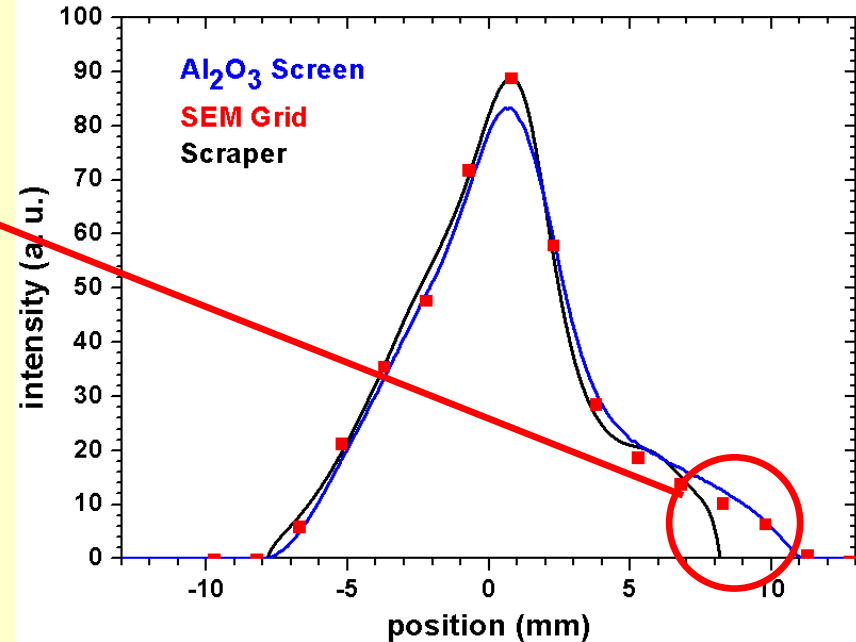
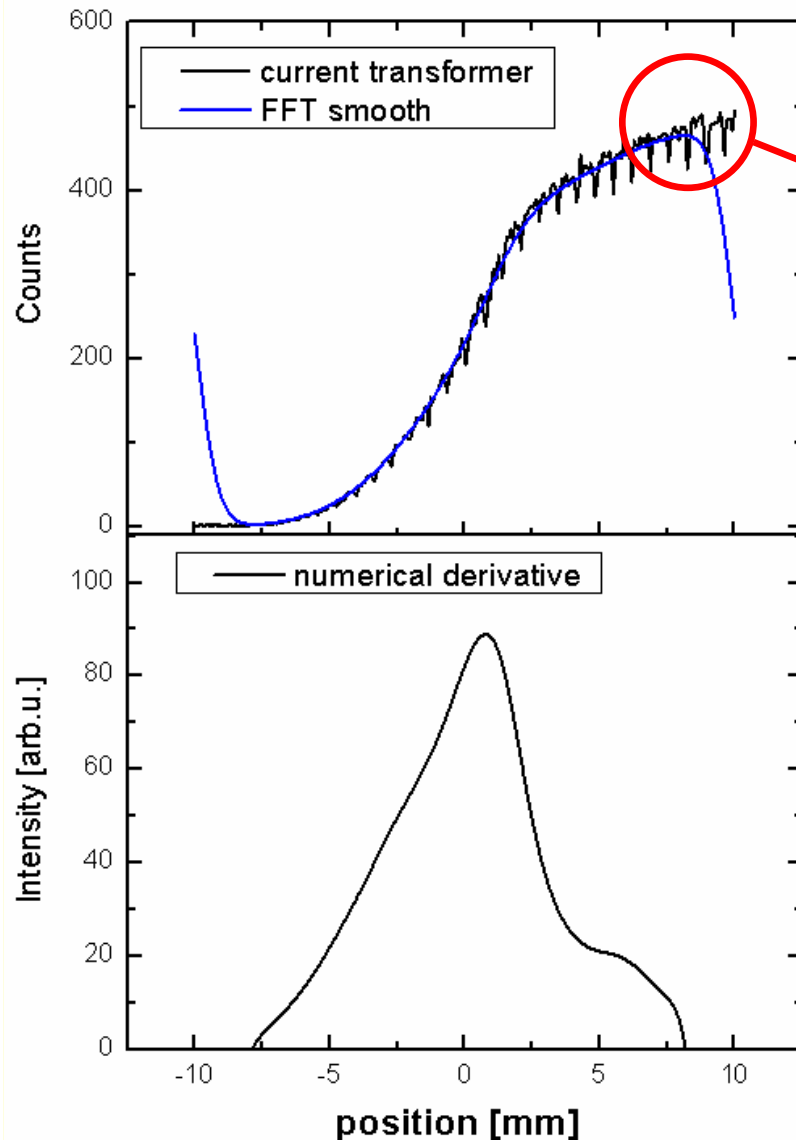
- How can one obtain a trusted beam profile with a better spatial resolution than a SEM-Grid

One can try to obtain a beam profile by using a scraper



SE suppression electronic ensures that no electrons affect the transformer measurements

Evaluation of data and comparison of diff. methods



Results:

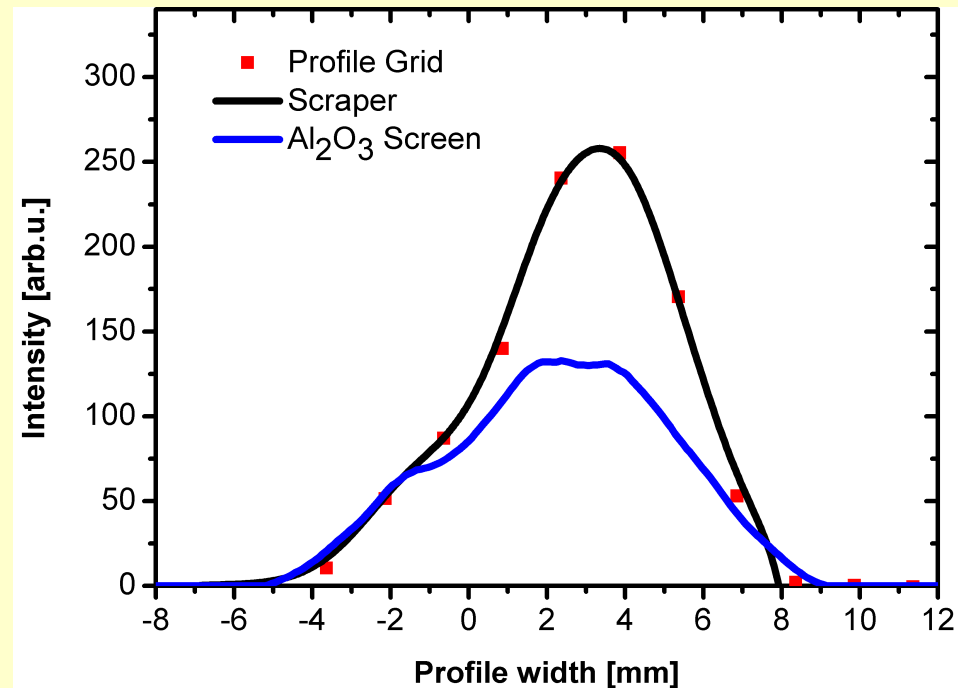
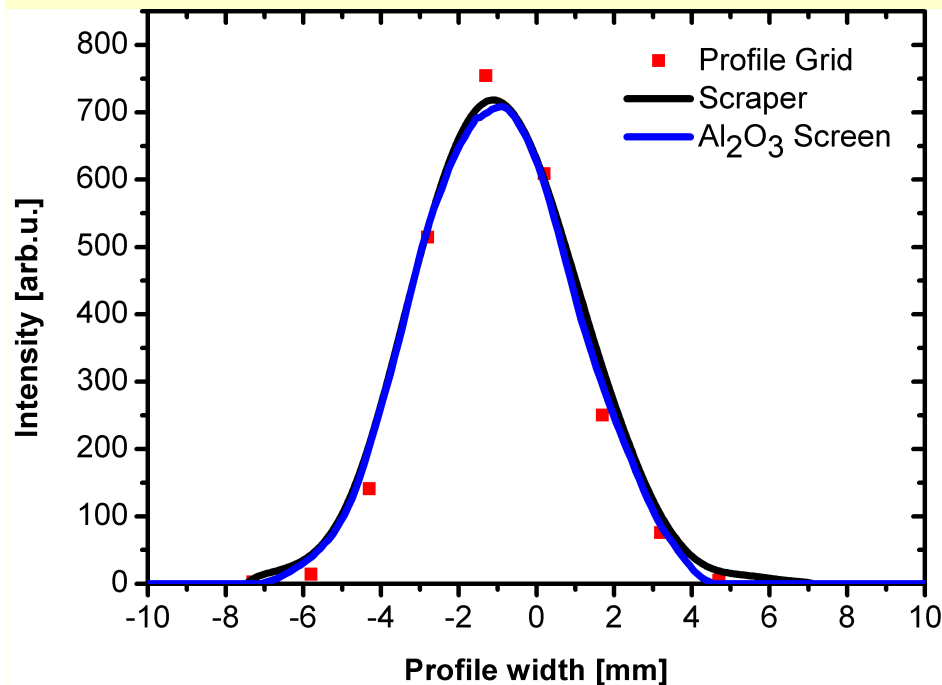
- SEM Grid and Scraper are in good agreement.
→ One can obtain a beam profile with a scraper with a much higher resolution than a SEM Grid
- **Allows to determine the response of the scintillator**
- Method needs a stable ion beam

What is the useful operating range of an Al_2O_3 Screen?

4.8 MeV/u
 $4.3 \cdot 10^{10}$ ppp
5ms

← Same pulse energy →

11.4 MeV/u
 $1.8 \cdot 10^{10}$ ppp
1,2 ms



Result: Light yield is the same for both energies. For the 11.4 MeV/u case, the imaged beam profile does not match to both reference methods. No $F^0 \rightarrow F^+$ conversion (Spectrum)

What could be an explanation for the results?

There are models that describe the light output of scintillators, but:

- for single particles
- only one species, e.g. Tl⁺
- can not predict changes in spectrum for diff. ions
- low doses (no damage)
-

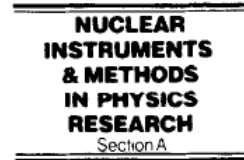
Light output is proportional to dose (e-h pairs) up to a quenching density ρ_q , above this dose the light output is constant. Fitting parameter: ρ_q

Different ansatz with δ electrons. 4 fitting parameters

Due to the complexity of the scintillating mechanisms, it is still under investigation



Nuclear Instruments and Methods in Physics Research A 356 (1995) 297–303



Scintillation response of nuclear particle detectors

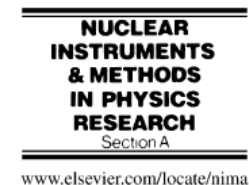
K. Michaelian, A. Menchaca-Rocha, E. Belmont-Moreno

Instituto de Física, Universidad Nacional Autónoma de México, A.P. 20-364, 01000 México D.F., Mexico

Received 26 July 1994



Nuclear Instruments and Methods in Physics Research A 482 (2002) 674–692

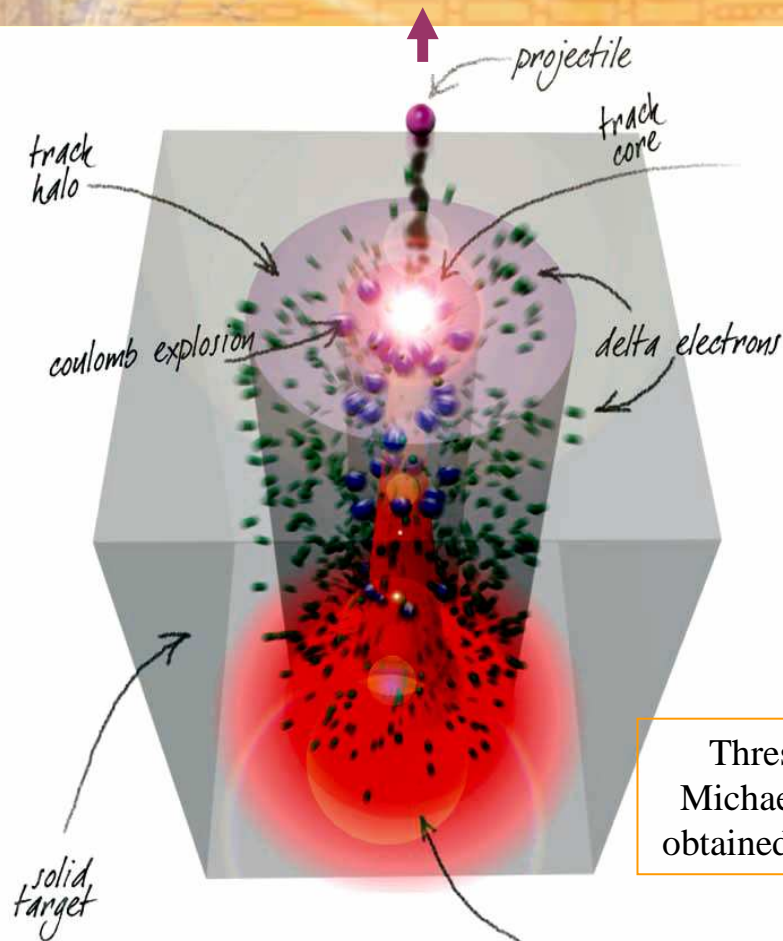


Response of CsI(Tl) scintillators over a large range in energy and atomic number of ions Part I: recombination and δ -electrons

M. Pârlog^{a,b}, B. Borderie^{b,*}, M.F. Rivet^b, G. Tăbăcaru^{a,b}, A. Chbihi^c,



The radial dose distribution of an ion track

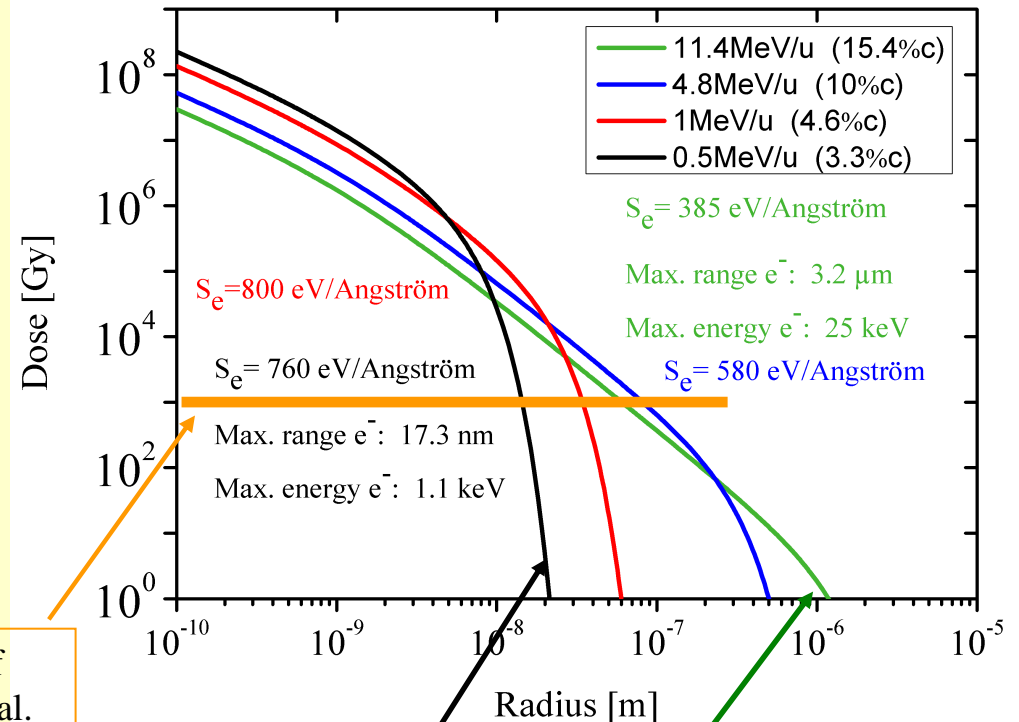


Courtesy of Marek Skupinski (Uppsala)

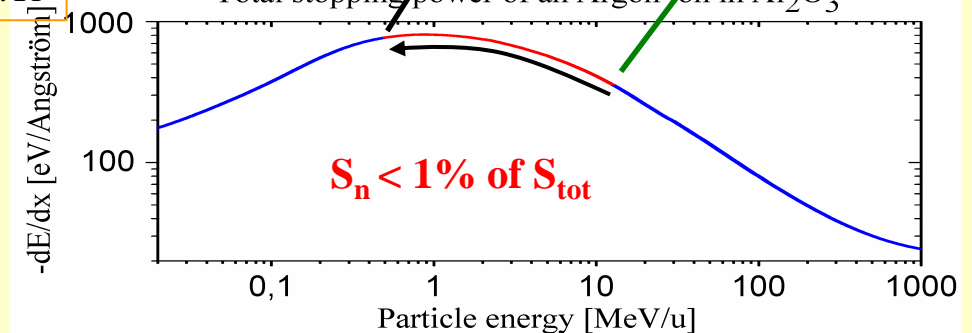
Threshold of Michaelian et al. obtained for CsI:TI

The camera can see all energies (depth) weighted with a Lambert-Beer absorption

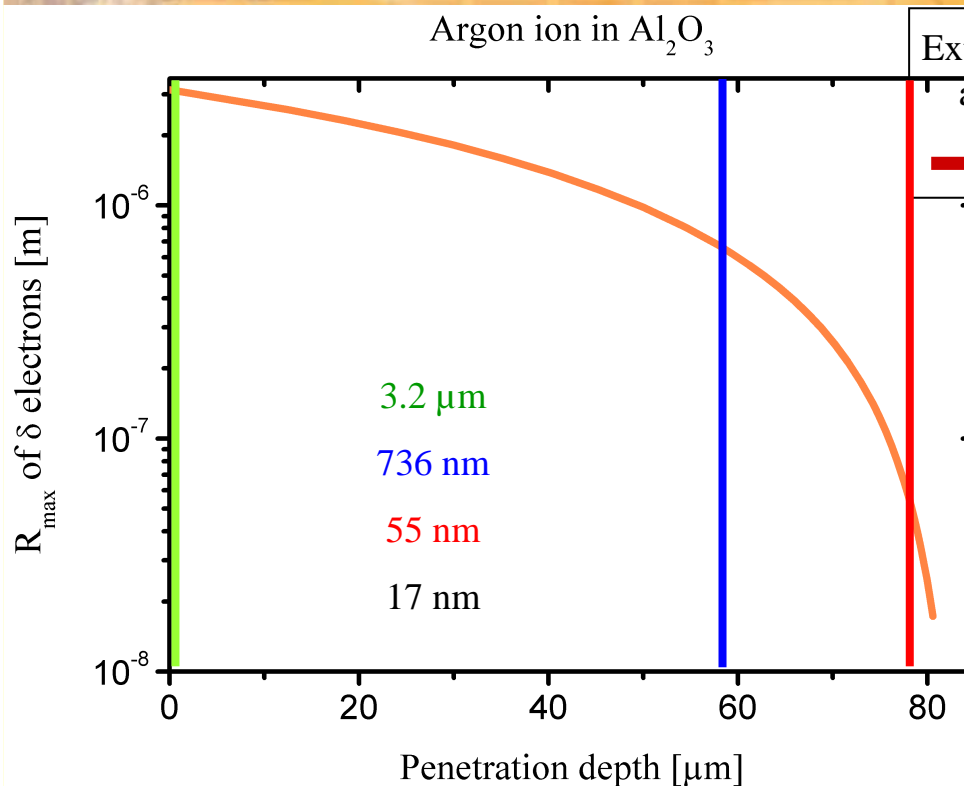
Radial dose distribution for an Argon ion in Al_2O_3



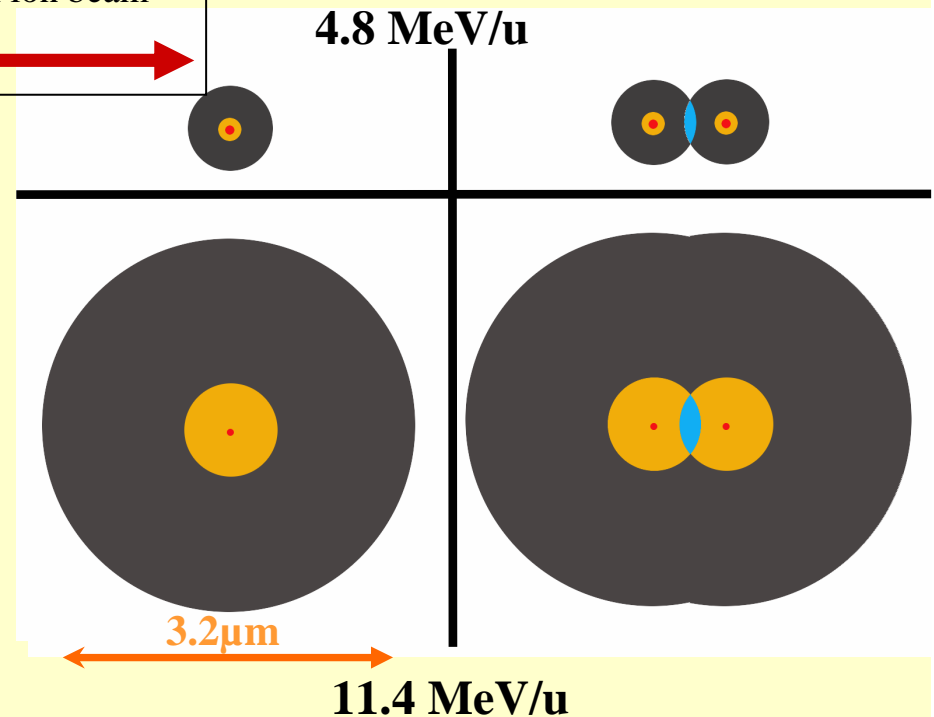
Total stopping power of an Argon ion in Al_2O_3



Penetration depth and R_{\max}



Schematic overlap of ion excitation tracks in space and time



Particle energy: 11.4 MeV/u, 4.8 MeV/u, 1 MeV/u

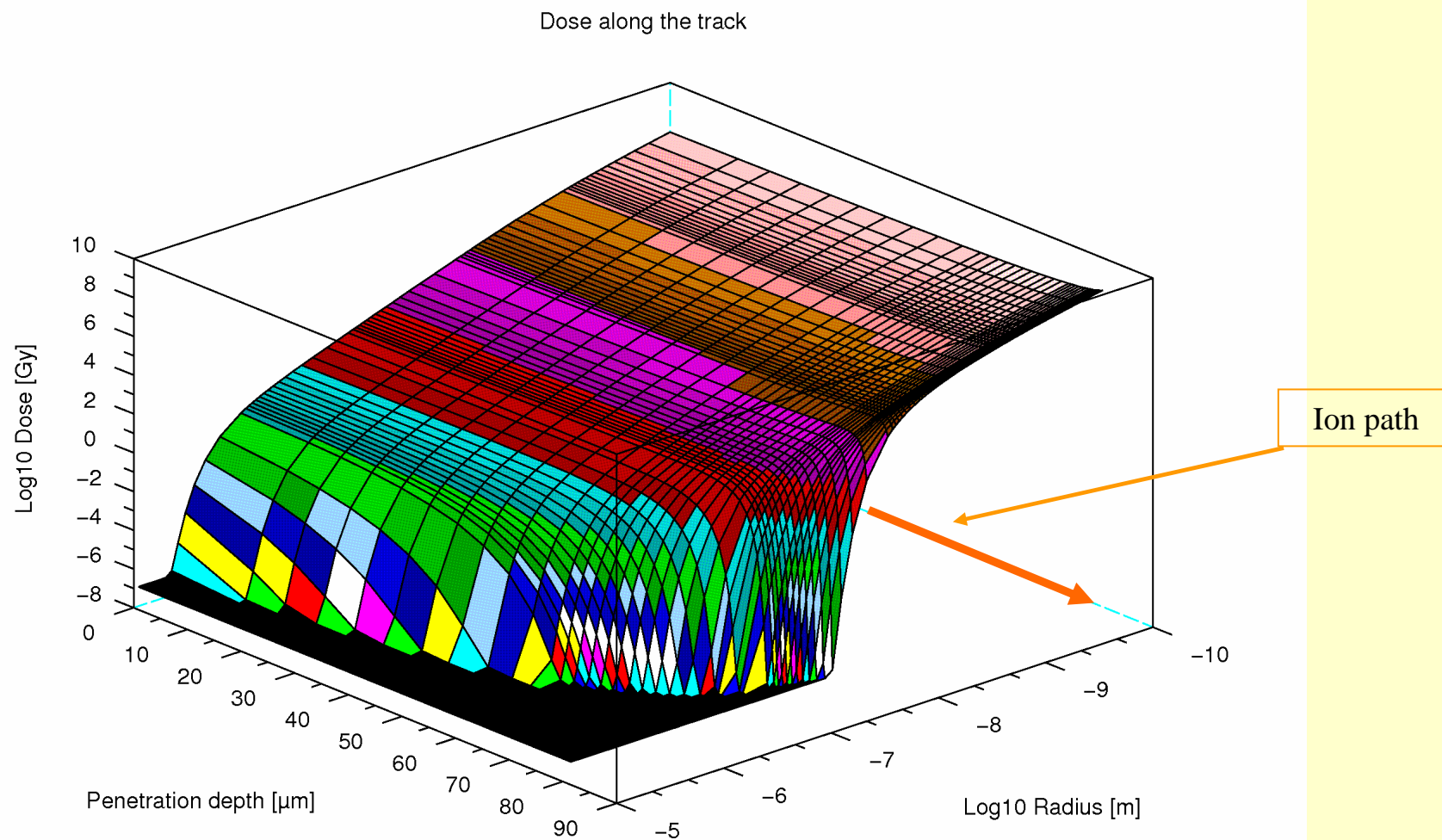
1E9 particles/cm² within the lifetime of the state.

The camera can see all energies (depth) weighted with a Lambert-Beer absorption

Result:

Higher light yield for faster particles but on the other hand, overlapping excitation tracks in the centre of the ion beam → F^0 centre could be ionized again! → → Spectrum?

Penetration depth and R_{\max}



The Model

For a time dependent 4D Monte Carlo Model, e.g. for Al_2O_3 , one would need in my opinion the cross-sections for:

- Electron capture at F^{2+} and F^+
- Hole capture at F^0 and F^+
- Ionization of F^+ and F^0

..... and the hole dynamic of charge carrier production, movement and trapping in the bulk material in dependence of ionization density.

Each one of them is a separate PhD-thesis, and it seem difficult to me to measure them independently of each other.

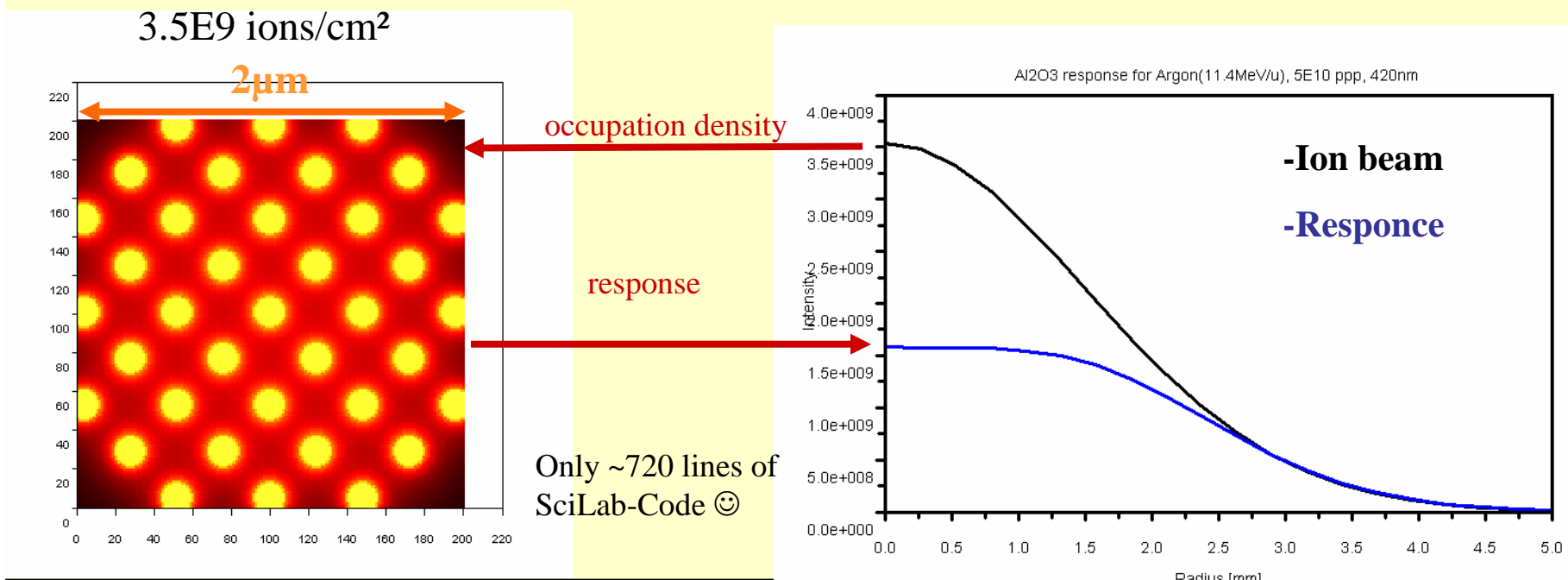
So lets try an time independent model

The Model

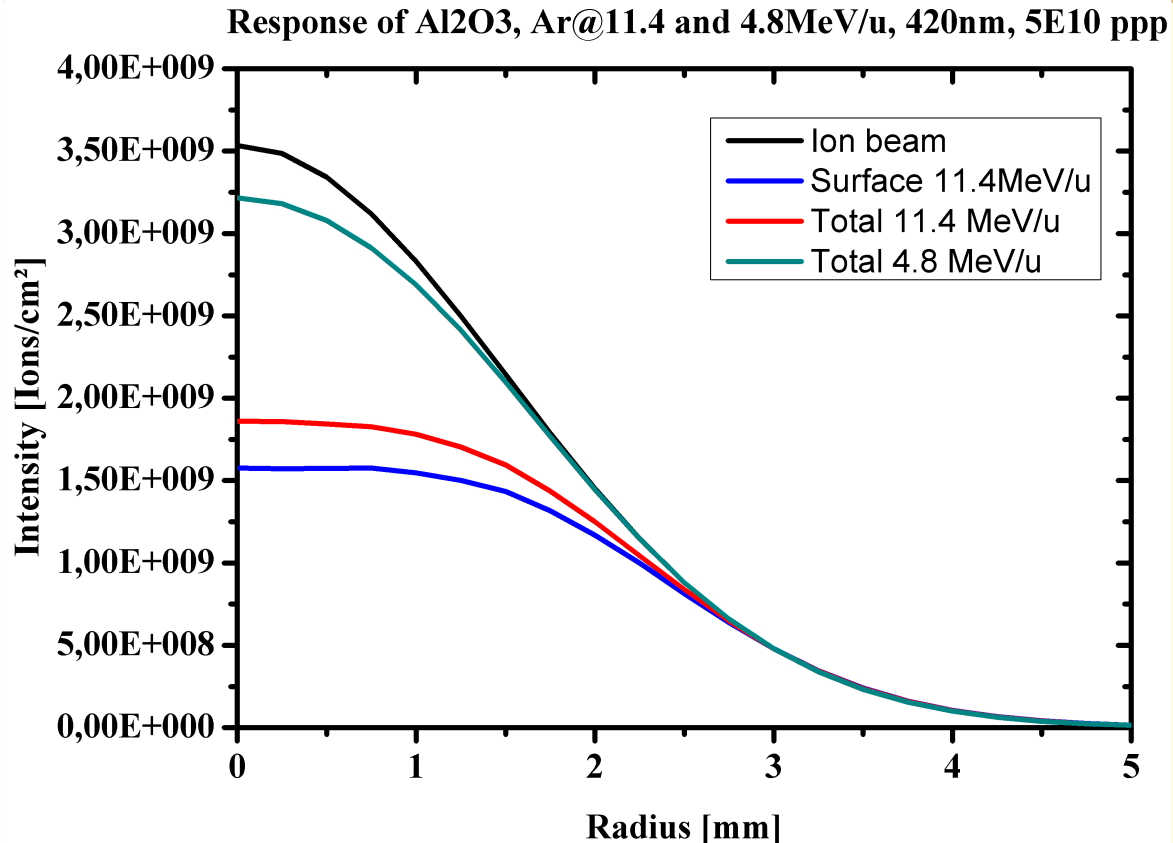
Assumption:

1. The radial dose distribution with the parametrization of Katz et al. '96 is valid
2. There is an ionization threshold like the one proposed by Michaelian et al.
3. The re-ionization process has a linear behaviour
4. The pulse length is smaller than e-h recombination+lifetime of the state (at least valid for F^{0*} state of Al_2O_3)

!The only fitting parameter is the ionization threshold ρ_i !



The prediction..... 'Tadaaaa'



5E10 ppp

ionization threshold = 5E2 Gy

Results: The Model fits the experimental results quite well

For a **gaussian ion beam**, Al₂O₃ screens, Argon @ 11.4 MeV/u, and 5E10 ppp;

The projection of a 4.8MeV/u ion beam is way less deformed then the one of 11.4 MeV/u, if one looks at the F⁰ (420nm) emission → F⁺ (330)nm. And there is no big contribution to the signal from the end of the ion track.

What about some other materials....

Herasil is not suitable for high current due to:

- Crack formation
- Has a threshold for light-output → measures wrong (The smallest beam profile is not always the correct one)
- Can have reflections from the back-side, due to its transparency
- Very low light-output

ZrO₂:Mg (Z507) is suitable for high current operation, but

- Has lower light-output than Al₂O₃
- “Saturates” earlier than Al₂O₃

ZrO₂:Y (Z700) is suitable for high current operation, but

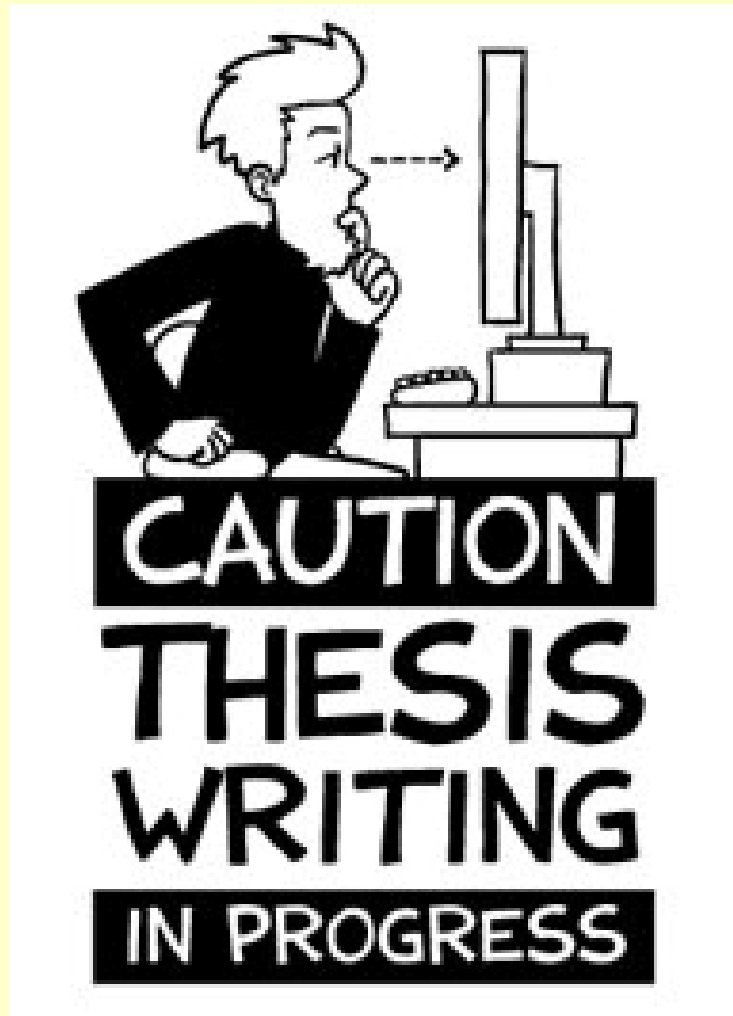
- Has a threshold for light-output → measures wrong
- Has low light-output

and the winner is: **Al₂O₃**

Summary of 2010



Winter / Spring 2011



Should be
ready in April
2011

Summer 2011



The defenses

2011, August





Fätsch